

Methodology Article

Solving Nonlinear Evolution Equations by (G'/G) -Expansion Method

Attia Rani¹, Munazza Saeed¹, Muhammad Ashraf¹, Rakshanda Zaman¹,
Qazi Mahmood-Ul-Hassan^{1,*}, Kamran Ayub², Muhammad Yaqub Khan², Madiha Afzal³

¹Department of Mathematics, University of Wah, Wah Cantt., Pakistan

²Department of Mathematics, Riphah International University, Islamabad, Pakistan

³Department of Mathematics, Allama Iqbal Open University, Islamabad, Pakistan

Email address:

qazimahmood@uow.edu.pk (Q. Mahmood-Ul-Hassan)

*Corresponding author

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Abstract: Nonlinear mathematical models and their solutions attain much attention in soliton theory. In this paper, main focus is to find travelling wave solutions of foam drainage equation and NLEE of fourth order. (G'/G) -expansion method is applied on these nonlinear differential equations. Wave transformation is used to convert nonlinear partial differential equation into an ordinary differential equation. It is observed that (G'/G) -expansion method is advanced and easy tool for finding solution of NLEEs in engineering, optics and mathematical physics. The proposed method is highly effective and reliable.

Keywords: (G'/G) -Expansion Method, Nonlinear Evolution Equations, Travelling Wave Solutions, Maple 18

1. Introduction

In the last few years, we have observed an extraordinary progress in soliton theory. Solitons have been studied by various mathematician, physicists and engineers for their applications in physical phenomena's. Firstly soliton waves are observed by an engineer John Scott Russell. Wide ranges of phenomena in mathematics and physics are modeled by differential equations. In nonlinear science it is of great importance and interest to explain physical models and attain analytical solutions. In the recent past large series of chemical, biological and physical singularities are feint by nonlinear partial differential equations. At present the prominent and valuable progress are made in the field of physical sciences. The great achievement is the development of various techniques to hunt for solitary wave solutions of differential equations. In nonlinear physical sciences, an essential contribution is of exact solutions because of this we can study physical behaviors and discuss more features of the problem

which give direction to more applications.

A reliable technique presented by Wang et al which is known as (G'/G) -expansion method for the nonlinear evolution equations (NLEEs) and provides the exact traveling wave solutions. In this technique, a linear ordinary differential equation of second order $G''(\eta) + \sigma G'(\eta) + \varphi G(\eta)$ is used, as the auxiliary equation. (G'/G) -expansion method [1-16] applied to solve the various types of the nonlinear evolution equations. A new modification introduced by Zhao et al. [17] in (G'/G) -expansion method. Extended (G'/G) -expansion method for solving (2+1)-dimensional NLEEs latter on introduce by Zayad and Abdelaziz [18]. For NLEEs arising in engineering and mathematical physics, multiple (G'/G) -expansion method applied by Chen and Li [19] in the recent year. (G'/G) -expansion method applied by Wang et al. [20] to find the traveling wave solutions of the Broer-Kaup and the approximate long water wave equations. For obtaining the traveling wave solutions of NLEEs (G'/G) -expansion method used by Liu et al and Aslan and Ozis [21] to find out the exact solutions of NLEEs. In the proposed method the solutions are expressed in term of rational trigonometric, the

rational hyperbolic and rational functions. The proposed method is powerful tool and very user-friendly for solving NLEEs. On exact solution some novel results and computational methods involved to travelling-wave transformation, see the references, [22-30].

2. Analysis of (G'/G) -Expansion Method

The general form of nonlinear partial differential equation is

$$\bar{P}(\bar{v}, \bar{v}_t, \bar{v}_x, \bar{v}_y, \bar{v}_z, \bar{v}_{tt}, \bar{v}_{xx}, \bar{v}_{yy}, \bar{v}_{zz}, \bar{v}_{xt}, \bar{v}_{yt}, \bar{v}_{zt}, \bar{v}_{xy}, \bar{v}_{xz}, \bar{v}_{yz}, \dots) = 0. \tag{1}$$

Here \bar{P} is a polynomial in $\bar{u}(x, t)$. The steps of (G'/G) -expansion method are as follows:

Step 1: Seek travelling wave variable of Eq. (1) by letting

$$\bar{v}(x, t) = \bar{v}(\bar{\xi}), \bar{\xi} = kx + ly + mz + \omega t.$$

and Eq. (1) transform into the ODE.

$$Q(\bar{v}, \bar{v}', \bar{v}'', \bar{v}''', \dots) = 0. \tag{2}$$

Here primes are representing the derivative of \bar{v} with respect to $\bar{\xi}$ and ω denotes constant.

Step 2: Constant (s) of integration can be obtained by integrating Eq. (2) term wise one or more times, if possible. For minimalism, the integration constant (s) can be set equal

$$\begin{aligned} \frac{G'(\bar{\xi})}{G(\bar{\xi})} &= \frac{\sqrt{\sigma^2 - 4\varphi}}{2} \left(\frac{c_1 \sinh\left(\frac{1}{2}\sqrt{\sigma^2 - 4\varphi}\bar{\xi}\right) + c_2 \cosh\left(\frac{1}{2}\sqrt{\sigma^2 - 4\varphi}\bar{\xi}\right)}{c_1 \cosh\left(\frac{1}{2}\sqrt{\sigma^2 - 4\varphi}\bar{\xi}\right) + c_2 \sinh\left(\frac{1}{2}\sqrt{\sigma^2 - 4\varphi}\bar{\xi}\right)} \right) - \frac{\sigma}{2}, \sigma^2 - 4\varphi > 0. \\ &= \frac{\sqrt{-\sigma^2 + 4\varphi}}{2} \left(\frac{-c_1 \sin\left(\frac{1}{2}\sqrt{-\sigma^2 + 4\varphi}\bar{\xi}\right) + c_2 \cos\left(\frac{1}{2}\sqrt{-\sigma^2 + 4\varphi}\bar{\xi}\right)}{c_1 \cos\left(\frac{1}{2}\sqrt{-\sigma^2 + 4\varphi}\bar{\xi}\right) + c_2 \sin\left(\frac{1}{2}\sqrt{-\sigma^2 + 4\varphi}\bar{\xi}\right)} \right) - \frac{\sigma}{2}, \sigma^2 - 4\varphi < 0. \\ &= \frac{2c_1}{c_1 + c_2\bar{\xi}} - \frac{\sigma}{2}, \sigma^2 - 4\varphi = 0. \end{aligned} \tag{5}$$

Where c_1, c_2 are unknown constants and we have

$$\begin{aligned} \left(\frac{G'(\bar{\xi})}{G(\bar{\xi})}\right)' &= -\left[\left(\frac{G'}{G}\right)^2 + \sigma\left(\frac{G'}{G}\right) + \varphi\right]. \\ \left(\frac{G'(\bar{\xi})}{G(\bar{\xi})}\right)'' &= \left[2\left(\frac{G'}{G}\right)^3 + 3\sigma\left(\frac{G'}{G}\right)^2 + (\sigma^2 + 2\varphi)\left(\frac{G'}{G}\right) + \varphi\lambda\right]. \\ &\vdots \end{aligned} \tag{4}$$

Where the primes represents the derivatives w.r.t $\bar{\xi}$. To find out \bar{v} explicitly, we follow these four steps.

Step 4: Substituting Eq. (3) and (4) into Eq. (2) after that, collect all terms with the same order of (G'/G) together, the left-hand side of Eq. (1) is converted into a polynomial in (G'/G) . Then by setting each coefficient equal to zero in this polynomial yields a set of algebraic equations for $k, l, m, \omega, \varepsilon$ and $a_n, n = 0, 1, \dots, M$.

Step 5: Solve this system of algebraic equations obtained for $k, l, m, \omega, \varepsilon$ and $a_n, n = 0, 1, \dots, M$ by using MAPLE 18.

Step 6: Use the results obtained from above steps to get a series of fundamental solutions $\bar{v}(\bar{\xi})$ of Eq. (2) that is depending on $\left(\frac{G'}{G}\right)$. Since the solutions of Eq. (3) will be well known for us, and then we can get the exact solutions of Eq.

to zero.

Step 3: According to the given proposed algorithm, suppose that the wave solution can be written as follows

$$\bar{v}(\xi) = a_0 + \sum_{n=1}^M a_n \left(\frac{G'}{G}\right)^n. \tag{3}$$

Where G is the solution of 1st order nonlinear equation in the following form:

$$G'' + \sigma G' + \varphi G = 0. \tag{4}$$

Where σ and φ are unknown constants. By using the general solution of Eq. (3), we get

3. Numerical Applications

We solve the following two problems to illustrate the implementation of the (G'/G) expansion method.

3.1. Foam Drainage Equation

Consider the Foam drainage equation [8]

$$\frac{\partial \eta}{\partial t} + \frac{\partial}{\partial x} \left(\eta^2 - \frac{\sqrt{\eta}}{2} \frac{\partial \eta}{\partial x} \right) = 0. \tag{6}$$

Here x and t scaled position and time coordinates, η represents the cross section of a channel formed at the point of

intersection of three films, mostly indicated as Plateau border (liquid filled channels). We focus on quantitative description of the coupling of drainage. Foam drainage is the flow of liquid from Plateau borders and the point where four channels meet between the bubbles, derived from capillarity and gravity. In foam stability foam drainage plays an important role. In fact, structure of the foam becomes fragile, when foam dries. [21]

Using the transformation as $\eta = \phi(\bar{\xi}), \bar{\xi} = \gamma(x + \omega t)$, where ω and γ are unknown constants, the given Eq. (6) is being transferred to

$$\gamma\omega \frac{\partial \phi}{\partial \bar{\xi}} + \gamma \frac{\partial}{\partial \bar{\xi}} \left(\phi^2 - \frac{\gamma}{2} \sqrt{\phi} \frac{\partial \phi}{\partial \bar{\xi}} \right) = 0. \tag{7}$$

Integrating Eq. (7) once with respect to η and setting integration constant equal to zero, we have

$$\gamma\omega\phi + \gamma \left(\phi^2 - \frac{\gamma}{2} \sqrt{\phi} \frac{\partial \phi}{\partial \bar{\xi}} \right) = 0. \tag{8}$$

By substituting $\phi(\bar{\xi}) = v^2(\bar{\xi})$, we have

$$\gamma\omega v^2 + \gamma \left(v^4 - \frac{\gamma}{2} v \cdot 2vv' \right) = 0.$$

$$\frac{G'}{G} = \begin{cases} \frac{\sqrt{\lambda^2 - 4\mu}}{2} \left(\frac{k_1 \sinh\left(\frac{\sqrt{\lambda^2 - 4\mu}\bar{\xi}}{2}\right) + k_2 \cosh\left(\frac{\sqrt{\lambda^2 - 4\mu}\bar{\xi}}{2}\right)}{k_1 \cosh\left(\frac{\sqrt{\lambda^2 - 4\mu}\bar{\xi}}{2}\right) + k_2 \sinh\left(\frac{\sqrt{\lambda^2 - 4\mu}\bar{\xi}}{2}\right)} \right) - \frac{\lambda}{2}, \lambda^2 - 4\mu > 0. \\ \frac{\sqrt{4\mu - \lambda^2}}{2} \left(\frac{-k_1 \sinh\left(\frac{\sqrt{4\mu - \lambda^2}\bar{\xi}}{2}\right) + k_2 \cosh\left(\frac{\sqrt{4\mu - \lambda^2}\bar{\xi}}{2}\right)}{k_1 \cosh\left(\frac{\sqrt{4\mu - \lambda^2}\bar{\xi}}{2}\right) + k_2 \sinh\left(\frac{\sqrt{4\mu - \lambda^2}\bar{\xi}}{2}\right)} \right) - \frac{\lambda}{2}, \lambda^2 - 4\mu < 0. \\ \frac{2k_1}{k_1 + k_2\bar{\xi}} - \frac{\lambda}{2}, \lambda^2 - 4\mu = 0. \end{cases}$$

Putting Eq. (10) into Eq. (9) and by collecting all terms with the same order of $\left(\frac{G'}{G}\right)$, we get a set of algebraic equations for ω, a_0 and a_1 as follows

$$\left(\frac{G'}{G}\right)^0 : \omega + a_0^2 + a_1\lambda\mu = 0,$$

$$\left(\frac{G'}{G}\right)^1 : 2a_0a_1 + a_1\lambda^2 = 0,$$

$$\left(\frac{G'}{G}\right)^2 : a_1^2 + a_1\lambda = 0.$$

Constants ω, a_0 and a_1 can be determined by using MAPLE 18, we have one solution set.

$$a_0 = -\frac{1}{2}\lambda^2, a_1 = -\lambda, \omega = -\frac{1}{4}\lambda^4 + \lambda^2\mu. \tag{12}$$

By using values of the above constants in Eq. (10), we get

$$v = -\lambda \left(\frac{G'}{G}\right) - \frac{1}{2}\lambda^2, \omega = -\frac{1}{4}\lambda^4 + \lambda^2\mu. \tag{13}$$

Or equivalently

$$\omega + v^2 - \gamma v' = 0. \tag{9}$$

By using the homogenous principle we balance the u' and v^2 , we get

$$M + 1 = 2M,$$

$$M = 1.$$

We assume equation Eq. (9) has the solution

$$v = a_0 + a_1 \left(\frac{G'}{G}\right), a_1 \neq 0. \tag{10}$$

Here a_0 and a_1 are unknown constants to be find out later. $G = G(\bar{\xi})$ satisfy the second order linear ordinary differential equation of the form

$$G''(\bar{\xi}) + \lambda G'(\bar{\xi}) + \mu G(\bar{\xi}) = 0. \tag{11}$$

λ and μ are constants, from Eq. (11), we have

Case I: When $\lambda^2 - 4\mu > 0$.

$$v_1 = -\frac{\lambda\sqrt{\lambda^2 - 4\mu}}{2} \left(\frac{k_1 \sinh\left(\frac{\sqrt{\lambda^2 - 4\mu}\bar{\xi}}{2}\right) + k_2 \cosh\left(\frac{\sqrt{\lambda^2 - 4\mu}\bar{\xi}}{2}\right)}{k_1 \cosh\left(\frac{\sqrt{\lambda^2 - 4\mu}\bar{\xi}}{2}\right) + k_2 \sinh\left(\frac{\sqrt{\lambda^2 - 4\mu}\bar{\xi}}{2}\right)} \right). \tag{14}$$

Here k_1 and k_2 are arbitrary constants.

If $k_1 = 0$, then solution (14) can be simplified as

$$v_2 = -\frac{\lambda\sqrt{\lambda^2 - 4\mu}}{2} \coth\left(\frac{\sqrt{\lambda^2 - 4\mu}\bar{\xi}}{2}\right). \tag{15}$$

If $k_2 = 0$, then solution (14) can be simplified as

$$v_3 = -\frac{\lambda\sqrt{\lambda^2 - 4\mu}}{2} \tanh\left(\frac{\sqrt{\lambda^2 - 4\mu}\bar{\xi}}{2}\right). \tag{16}$$

Case II: When $\lambda^2 - 4\mu < 0$.

$$v_4 = -\frac{\lambda\sqrt{4\mu-\lambda^2}}{2} \left(\frac{-k_1 \sinh\left(\frac{\sqrt{4\mu-\lambda^2}\xi}{2}\right) + k_2 \cosh\left(\frac{\sqrt{4\mu-\lambda^2}\xi}{2}\right)}{k_1 \cosh\left(\frac{\sqrt{4\mu-\lambda^2}\xi}{2}\right) + k_2 \sinh\left(\frac{\sqrt{4\mu-\lambda^2}\xi}{2}\right)} \right). \quad (17)$$

If $k_1 = 0$, then solution (17) can be simplified as

$$v_5 = -\frac{\lambda\sqrt{4\mu-\lambda^2}}{2} \coth\left(\frac{\sqrt{4\mu-\lambda^2}\xi}{2}\right). \quad (18)$$

If $k_2 = 0$, then solution (17) can be simplified as

$$v_6 = \frac{\lambda\sqrt{4\mu-\lambda^2}}{2} \tanh\left(\frac{\sqrt{4\mu-\lambda^2}\xi}{2}\right). \quad (19)$$

Case III: When $\lambda^2 - 4\mu = 0$.

$$v_7 = -\frac{2\lambda k_1}{k_1 + k_2 \xi}. \quad (20)$$

Here in all the cases $\xi = \gamma \left(x + \left(-\frac{1}{4}\lambda^4 + \lambda^2\mu \right) t \right)$.

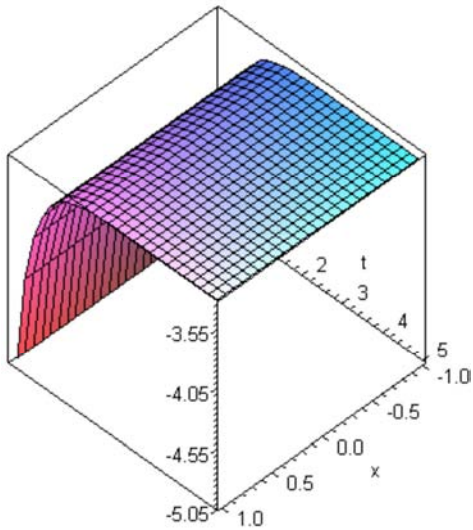


Figure 1. Soliton solution of $v_1(x,t)$ for $k_1 = 1, k_2 = 1.5, \gamma = 1, \lambda = 3, \mu = 1$.

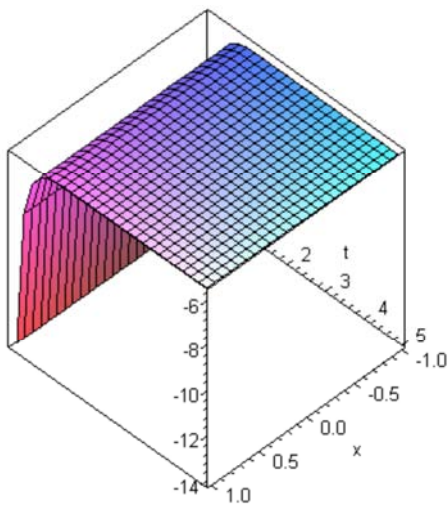


Figure 2. Soliton solution of $v_1(x,t)$ for $k_1 = 1, k_2 = 2.5, \gamma = 1.5, \lambda = 4, \mu = 2$.

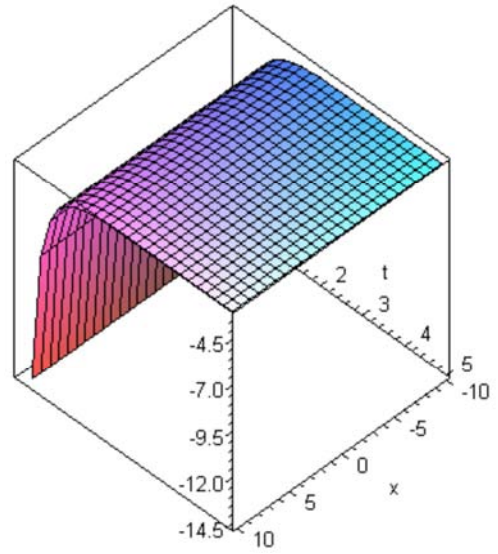


Figure 3. Soliton solution of $v_2(x,t)$ for $\gamma = 1, \lambda = 3, \mu = 1$.

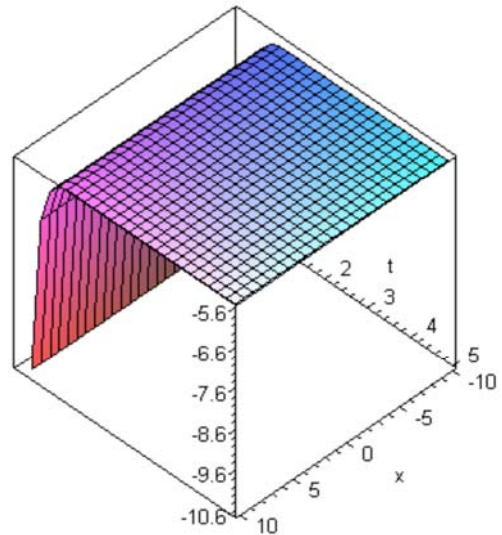


Figure 4. Soliton solution of $v_2(x,t)$ for $\gamma = 2, \lambda = 4, \mu = 2$.

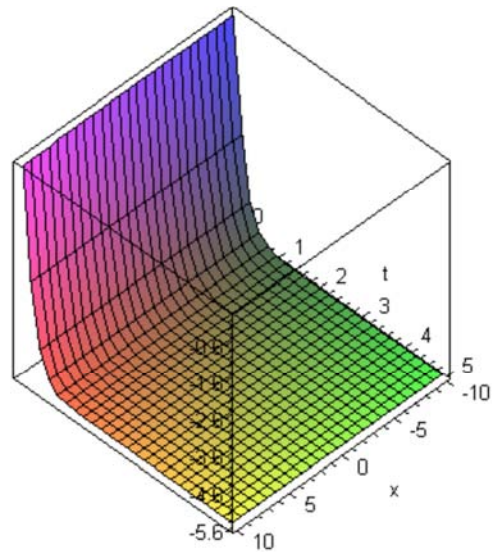


Figure 5. Soliton solution of $v_3(x,t)$ for $\gamma = 2, \lambda = 4, \mu = 2$.

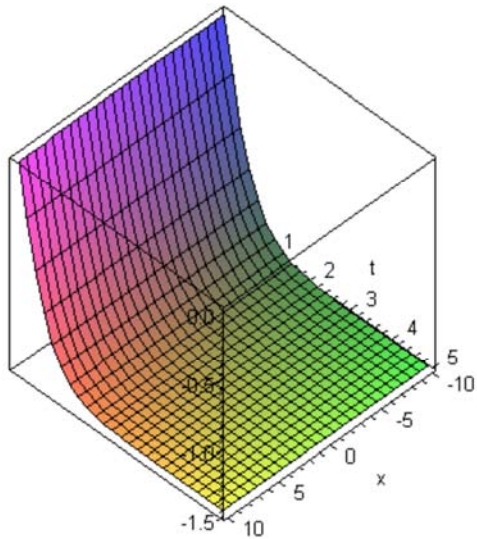


Figure 6. Soliton solution of $v_3(x,t)$ for $\gamma = 2.5, \lambda = 3, \mu = 2$.

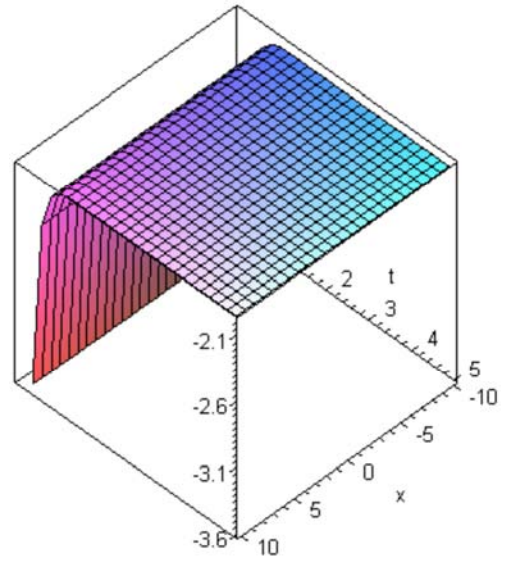


Figure 9. Soliton solution of $v_5(x,t)$ for $\gamma = 3, \lambda = 2, \mu = 2$.

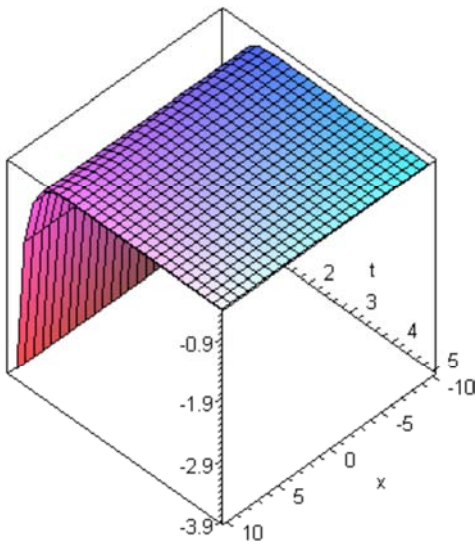


Figure 7. Soliton solution of $v_4(x,t)$ for $k_1 = 1, k_2 = 1.5, \gamma = 2.5, \lambda = 3, \mu = 3$.

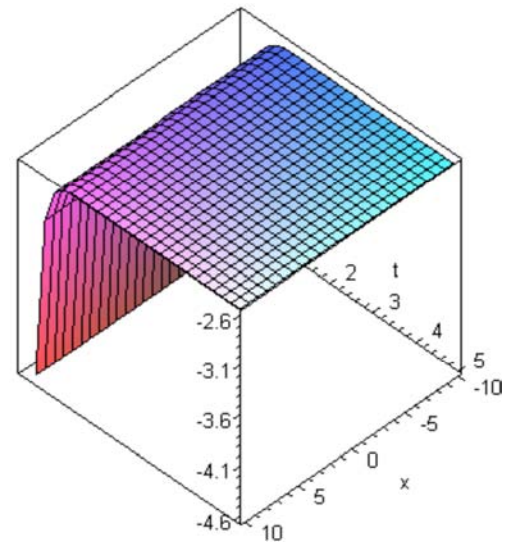


Figure 10. Soliton solution of $v_5(x,t)$ for $\gamma = 3.5, \lambda = 3, \mu = 3$.

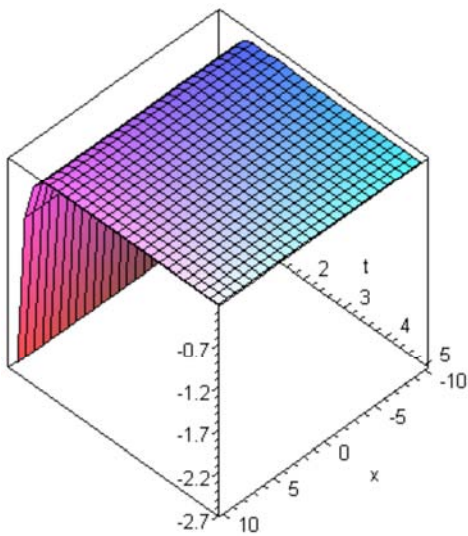


Figure 8. Soliton solution of $v_4(x,t)$ for $k_1 = 1.5, k_2 = 2, \gamma = 3, \lambda = 2, \mu = 2$.

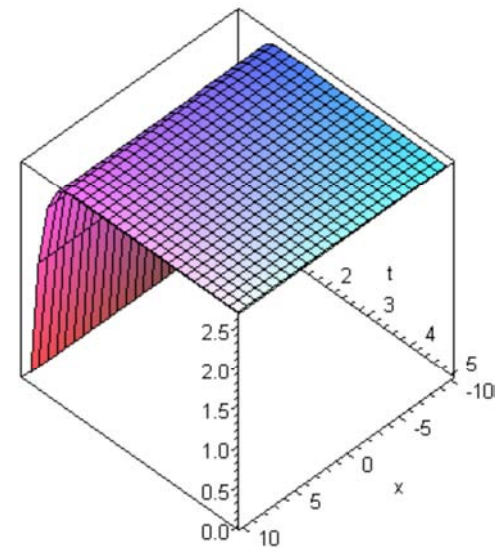


Figure 11. Soliton solution of $v_6(x,t)$ for $\gamma = 3.5, \lambda = 3, \mu = 3$.

3.2. Nonlinear Fourth Order Evolution Equation

Let us consider the NLEE of fourth order [13]

$$\eta_{tt} - a\eta_{xt}\eta_{xx} + b\eta_{xxxt} = 0. \tag{21}$$

Where a and b are constants. Fourth order NLEE is one of good starting point for study of non-linear water waves that was first point out by Dysthe in 1979. He obtained for gravity waves propagating at the interface of two superposed fluids of infinite depth over water in the presence of air flowing and a basic current shear.

By using the transformation as $\bar{\xi} = x - \omega t$, we reduce the given Eq. (21) into an ODE

$$\omega\eta'' - a(\eta'')^2 - b\eta^{(iv)} = 0. \tag{22}$$

By putting $u = \eta''$, we have

$$\omega u - au^2 - bu'' = 0. \tag{23}$$

Balancing the u'' and u^2 by using the homogenous principle, we get

$$M + 2 = 2M,$$

$$M = 2.$$

Now, we suppose Eq. (23) has the solution

$$u = a_0 + a_1 \left(\frac{G'}{G}\right) + a_2 \left(\frac{G'}{G}\right)^2, a_2 \neq 0. \tag{24}$$

Here a_0, a_1 and a_2 are constants to be find out later.

$G = G(\bar{\xi})$ satisfy the second order linear ordinary differential equation in the following form

$$G''(\bar{\xi}) + \lambda G'(\bar{\xi}) + \mu G(\bar{\xi}) = 0. \tag{25}$$

λ and μ are constants, from Eq. (25), we get

$$\frac{G'}{G} = \begin{cases} \frac{\sqrt{\lambda^2 - 4\mu}}{2} \left(\frac{k_1 \sinh\left(\frac{\sqrt{\lambda^2 - 4\mu}\bar{\xi}}{2}\right) + k_2 \cosh\left(\frac{\sqrt{\lambda^2 - 4\mu}\bar{\xi}}{2}\right)}{k_1 \cosh\left(\frac{\sqrt{\lambda^2 - 4\mu}\bar{\xi}}{2}\right) + k_2 \sinh\left(\frac{\sqrt{\lambda^2 - 4\mu}\bar{\xi}}{2}\right)} \right) - \frac{\lambda}{2}, \lambda^2 - 4\mu > 0. \\ \frac{\sqrt{4\mu - \lambda^2}}{2} \left(\frac{-k_1 \sinh\left(\frac{\sqrt{4\mu - \lambda^2}\bar{\xi}}{2}\right) + k_2 \cosh\left(\frac{\sqrt{4\mu - \lambda^2}\bar{\xi}}{2}\right)}{k_1 \cosh\left(\frac{\sqrt{4\mu - \lambda^2}\bar{\xi}}{2}\right) + k_2 \sinh\left(\frac{\sqrt{4\mu - \lambda^2}\bar{\xi}}{2}\right)} \right) - \frac{\lambda}{2}, \lambda^2 - 4\mu < 0. \\ \frac{2k_1}{k_1 + k_2\bar{\xi}} - \frac{\lambda}{2}, \lambda^2 - 4\mu = 0. \end{cases}$$

By putting the Eq. (24) into eq. (23) and collecting all terms with the same order of $\left(\frac{G'}{G}\right)$ we yields a set of algebraic equations for ω, a_0, a_1 and a_2 as follows

$$\left(\frac{G'}{G}\right)^0 : \omega a_0 - aa_0^2 - ba_1\lambda\mu - 2ba_2\mu^2 = 0,$$

$$\left(\frac{G'}{G}\right)^1 : \omega a_1 - 2aa_0a_1 - ba_1\lambda^2 - 2ba_1\mu - 6ba_2\lambda\mu = 0,$$

$$\left(\frac{G'}{G}\right)^2 : \omega a_2 - 2aa_0a_2 - aa_1^2 - 3ba_1\lambda - 4ba_2\lambda^2 - 8ba_2\mu = 0,$$

$$\left(\frac{G'}{G}\right)^3 : -2aa_1a_2 - 2ba_1 - 10ba_2\lambda = 0,$$

$$\left(\frac{G'}{G}\right)^4 : -aa_2^2 - 6ba_2 = 0.$$

Constants ω, a_0, a_1 and a_2 can be determined by using MAPLE 18, we have following two solution sets

1st Solution Set:

$$a_2 = -\frac{6b}{a}, a_1 = -\frac{6b\lambda}{a}, a_0 = -\frac{6b\mu}{a}, \omega = -4b\mu + b\lambda^2. \tag{26}$$

By substituting values of the above constants in Eq. (24), we have

$$u = -\frac{6b}{a} \left(\frac{G'}{G}\right)^2 - \frac{6b\lambda}{a} \left(\frac{G'}{G}\right) - \frac{6b\mu}{a}.$$

Case I: When $\lambda^2 - 4\mu > 0$.

$$u_1 = -3b \left(\frac{\lambda^2 - 4\mu}{2a}\right) \left(\frac{k_1 \sinh\left(\frac{\sqrt{\lambda^2 - 4\mu}\xi}{2}\right) + k_2 \cosh\left(\frac{\sqrt{\lambda^2 - 4\mu}\xi}{2}\right)}{k_1 \cosh\left(\frac{\sqrt{\lambda^2 - 4\mu}\xi}{2}\right) + k_2 \sinh\left(\frac{\sqrt{\lambda^2 - 4\mu}\xi}{2}\right)} \right)^2 - \frac{6b\mu}{a}. \tag{27}$$

$$\eta_1 = -3b \left(\frac{\lambda^2 - 4\mu}{2a}\right) \iint_0^{\bar{\xi}} \left(\frac{k_1 \sinh\left(\frac{\sqrt{\lambda^2 - 4\mu}\xi}{2}\right) + k_2 \cosh\left(\frac{\sqrt{\lambda^2 - 4\mu}\xi}{2}\right)}{k_1 \cosh\left(\frac{\sqrt{\lambda^2 - 4\mu}\xi}{2}\right) + k_2 \sinh\left(\frac{\sqrt{\lambda^2 - 4\mu}\xi}{2}\right)} \right)^2 d\bar{\xi} d\bar{\xi} - \frac{3b\mu\bar{\xi}^2}{a}. \tag{28}$$

Where k_1 and k_2 are arbitrary constants.

If $k_1 = 0$, then solution (27) and (28) can be expressed as

$$u_2 = -3b \left(\frac{\lambda^2 - 4\mu}{2a}\right) \coth^2\left(\frac{\sqrt{\lambda^2 - 4\mu}\xi}{2}\right) - \frac{6b\mu}{a}. \tag{29}$$

$$\eta_2 = -3b \left(\frac{\lambda^2 - 4\mu}{2a}\right) \iint_0^{\bar{\xi}} \coth^2\left(\frac{\sqrt{\lambda^2 - 4\mu}\xi}{2}\right) d\bar{\xi} d\bar{\xi} - \frac{3b\mu\bar{\xi}^2}{a}. \tag{30}$$

If $k_2 = 0$, then solution (27) and (28) can be expressed as

$$u_3 = -3b \left(\frac{\lambda^2 - 4\mu}{2a}\right) \tanh^2\left(\frac{\sqrt{\lambda^2 - 4\mu}\xi}{2}\right) - \frac{6b\mu}{a}. \tag{31}$$

$$\eta_3 = -3b \left(\frac{\lambda^2 - 4\mu}{2a}\right) \iint_0^{\bar{\xi}} \tanh^2\left(\frac{\sqrt{\lambda^2 - 4\mu}\xi}{2}\right) d\bar{\xi} d\bar{\xi} - \frac{3b\mu\bar{\xi}^2}{a}. \tag{32}$$

Case II: When $\lambda^2 - 4\mu < 0$.

$$u_4 = -3b \left(\frac{4\mu - \lambda^2}{2a}\right) \left(\frac{-k_1 \sinh\left(\frac{\sqrt{4\mu - \lambda^2}\xi}{2}\right) + k_2 \cosh\left(\frac{\sqrt{4\mu - \lambda^2}\xi}{2}\right)}{k_1 \cosh\left(\frac{\sqrt{4\mu - \lambda^2}\xi}{2}\right) + k_2 \sinh\left(\frac{\sqrt{4\mu - \lambda^2}\xi}{2}\right)} \right)^2 - \frac{6b\mu^2}{a}. \tag{33}$$

$$\eta_4 = -3b \left(\frac{4\mu - \lambda^2}{2a}\right) \iint_0^{\bar{\xi}} \left(\frac{k_1 \sinh\left(\frac{\sqrt{4\mu - \lambda^2}\xi}{2}\right) + k_2 \cosh\left(\frac{\sqrt{4\mu - \lambda^2}\xi}{2}\right)}{k_1 \cosh\left(\frac{\sqrt{4\mu - \lambda^2}\xi}{2}\right) + k_2 \sinh\left(\frac{\sqrt{4\mu - \lambda^2}\xi}{2}\right)} \right)^2 d\bar{\xi} d\bar{\xi} - \frac{3b\mu\bar{\xi}^2}{a}. \tag{34}$$

If $k_1 = 0$, then solution (33) and (34) can be simplified as

$$u_5 = -3b \left(\frac{4\mu - \lambda^2}{2a}\right) \coth^2\left(\frac{\sqrt{4\mu - \lambda^2}\xi}{2}\right) - \frac{6b\mu}{a}. \tag{35}$$

$$v_5 = -3b \left(\frac{4\mu - \lambda^2}{2a}\right) \iint_0^{\bar{\xi}} \coth^2\left(\frac{\sqrt{4\mu - \lambda^2}\xi}{2}\right) d\bar{\xi} d\bar{\xi} - \frac{3b\mu\bar{\xi}^2}{a}. \tag{36}$$

If $k_2 = 0$, then solution (33) and (34) can be simplified as

$$u_6 = -3b \left(\frac{4\mu - \lambda^2}{2a}\right) \tanh^2\left(\frac{\sqrt{4\mu - \lambda^2}\xi}{2}\right) - \frac{6b\mu}{a}. \tag{37}$$

$$\eta_5 = -3b \left(\frac{4\mu - \lambda^2}{2a}\right) \iint_0^{\bar{\xi}} \tanh^2\left(\frac{\sqrt{4\mu - \lambda^2}\xi}{2}\right) d\bar{\xi} d\bar{\xi} - \frac{3b\mu\bar{\xi}^2}{a}. \tag{38}$$

Case III: When $\lambda^2 - 4\mu = 0$.

$$u_7 = -24b \frac{k_1^2}{a(k_1 + k_2\bar{\xi})^2} - \frac{6b\mu}{a}. \tag{39}$$

$$\eta_7 = -24b \iint_0^{\bar{\xi}} \frac{k_1^2}{a(k_1 + k_2\bar{\xi})^2} d\bar{\xi} d\bar{\xi} - \frac{3b\mu\bar{\xi}^2}{a}. \tag{40}$$

Here in all the cases $\bar{\xi} = x - (-4b\mu + b\lambda^2)t$.

2nd Solution Set:

$$\left. \begin{aligned} a_0 &= -\frac{b(\lambda^2+2\mu)}{a}, a_1 = -\frac{6b\lambda}{a}, \\ a_2 &= \frac{-6b}{a}, \omega = -b\lambda^2 + 4b\mu. \end{aligned} \right\} \quad (41)$$

$$u = \frac{-6b}{a} \left(\frac{G'}{G}\right)^2 - \frac{6b\lambda}{a} \left(\frac{G'}{G}\right) - \frac{b(\lambda^2 + 2\mu)}{a}.$$

Case I: When $\lambda^2 - 4\mu > 0$,

By substituting values of above constants in Eq. (24), we have

$$u_8 = -3b \left(\frac{\lambda^2-4\mu}{2a}\right) \left(\frac{k_1 \sinh\left(\frac{\sqrt{\lambda^2-4\mu}\xi}{2}\right) + k_2 \cosh\left(\frac{\sqrt{\lambda^2-4\mu}\xi}{2}\right)}{k_1 \cosh\left(\frac{\sqrt{\lambda^2-4\mu}\xi}{2}\right) + k_2 \sinh\left(\frac{\sqrt{\lambda^2-4\mu}\xi}{2}\right)} \right)^2 - \frac{b(\lambda^2+2\mu)}{a}. \quad (42)$$

$$\eta_8 = -3b \left(\frac{\lambda^2-4\mu}{2a}\right) \iint_0^{\bar{\xi}} \left(\frac{k_1 \sinh\left(\frac{\sqrt{\lambda^2-4\mu}\xi}{2}\right) + k_2 \cosh\left(\frac{\sqrt{\lambda^2-4\mu}\xi}{2}\right)}{k_1 \cosh\left(\frac{\sqrt{\lambda^2-4\mu}\xi}{2}\right) + k_2 \sinh\left(\frac{\sqrt{\lambda^2-4\mu}\xi}{2}\right)} \right)^2 d\bar{\xi} d\bar{\xi} - \frac{b(\lambda^2+2\mu)\bar{\xi}^2}{2a}. \quad (43)$$

Where k_1 and k_2 are arbitrary constants.

If $k_1 = 0$, then solution (42) and (43) can be simplified as

$$u_9 = -3b \left(\frac{\lambda^2-4\mu}{2a}\right) \coth^2 \left(\frac{\sqrt{\lambda^2-4\mu}\xi}{2}\right) - \frac{b(\lambda^2+2\mu)}{a}. \quad (44)$$

$$\eta_9 = -3b \left(\frac{\lambda^2-4\mu}{2a}\right) \iint_0^{\bar{\xi}} \coth^2 \left(\frac{\sqrt{\lambda^2-4\mu}\xi}{2}\right) d\bar{\xi} d\bar{\xi} - \frac{b(\lambda^2+2\mu)\bar{\xi}^2}{2a}. \quad (45)$$

If $k_2 = 0$, then solution (42) and (43) can be simplified as

$$u_{10} = -3b \left(\frac{\lambda^2-4\mu}{2a}\right) \tanh^2 \left(\frac{\sqrt{\lambda^2-4\mu}\xi}{2}\right) - \frac{b(\lambda^2+2\mu)}{a}. \quad (46)$$

$$\eta_{10} = -3b \left(\frac{\lambda^2-4\mu}{2a}\right) \iint_0^{\bar{\xi}} \tanh^2 \left(\frac{\sqrt{\lambda^2-4\mu}\xi}{2}\right) d\bar{\xi} d\bar{\xi} - \frac{b(\lambda^2+2\mu)\bar{\xi}^2}{2a}. \quad (47)$$

Case II: When $\lambda^2 - 4\mu < 0$,

$$u_{11} = -3b \left(\frac{4\mu-\lambda^2}{2a}\right) \left(\frac{-k_1 \sinh\left(\frac{\sqrt{4\mu-\lambda^2}\xi}{2}\right) + k_2 \cosh\left(\frac{\sqrt{4\mu-\lambda^2}\xi}{2}\right)}{k_1 \cosh\left(\frac{\sqrt{4\mu-\lambda^2}\xi}{2}\right) + k_2 \sinh\left(\frac{\sqrt{4\mu-\lambda^2}\xi}{2}\right)} \right)^2 - \frac{b(\lambda^2+2\mu)}{a}. \quad (48)$$

$$\eta_{11} = -3b \left(\frac{4\mu-\lambda^2}{2a}\right) \iint_0^{\bar{\xi}} \left(\frac{k_1 \sinh\left(\frac{\sqrt{4\mu-\lambda^2}\xi}{2}\right) + k_2 \cosh\left(\frac{\sqrt{4\mu-\lambda^2}\xi}{2}\right)}{k_1 \cosh\left(\frac{\sqrt{4\mu-\lambda^2}\xi}{2}\right) + k_2 \sinh\left(\frac{\sqrt{4\mu-\lambda^2}\xi}{2}\right)} \right)^2 d\bar{\xi} d\bar{\xi} - \frac{b(\lambda^2+2\mu)\bar{\xi}^2}{2a}. \quad (49)$$

If $k_1 = 0$, then solution (48) and (49) can be simplified as

$$u_{12} = -3b \left(\frac{4\mu-\lambda^2}{2a}\right) \coth^2 \left(\frac{\sqrt{4\mu-\lambda^2}\xi}{2}\right) - \frac{b(\lambda^2+2\mu)}{a}. \quad (50)$$

$$\eta_{12} = -3b \left(\frac{4\mu-\lambda^2}{2a}\right) \iint_0^{\bar{\xi}} \coth^2 \left(\frac{\sqrt{4\mu-\lambda^2}\xi}{2}\right) d\bar{\xi} d\bar{\xi} - \frac{b(\lambda^2+2\mu)\bar{\xi}^2}{2a}. \quad (51)$$

If $k_2 = 0$, then solution (48) and (49) can be simplified as

$$u_{13} = -3b \left(\frac{4\mu - \lambda^2}{2a} \right) \tanh^2 \left(\frac{\sqrt{4\mu - \lambda^2} \bar{\xi}}{2} \right) - \frac{b(\lambda^2 + 2\mu)}{a}. \tag{52}$$

$$\eta_{13} = -3b \left(\frac{4\mu - \lambda^2}{2a} \right) \iint_0^{\bar{\xi}} \tanh^2 \left(\frac{\sqrt{4\mu - \lambda^2} \bar{\xi}}{2} \right) d\bar{\xi} d\bar{\xi} - \frac{b(\lambda^2 + 2\mu)\bar{\xi}^2}{2a}. \tag{53}$$

Case III: When $\lambda^2 - 4\mu = 0$,

$$u_{14} = -24b \frac{k_1^2}{a(k_1 + k_2 \bar{\xi})^2} - \frac{b(\lambda^2 + 2\mu)}{a}. \tag{54}$$

$$\eta_{14} = -24b \iint_0^{\bar{\xi}} \frac{k_1^2}{a(k_1 + k_2 \bar{\xi})^2} d\bar{\xi} d\bar{\xi} - \frac{b(\lambda^2 + 2\mu)\bar{\xi}^2}{2a}. \tag{55}$$

Here in all the cases $\bar{\xi} = x - (4b\mu - b\lambda^2)t$.

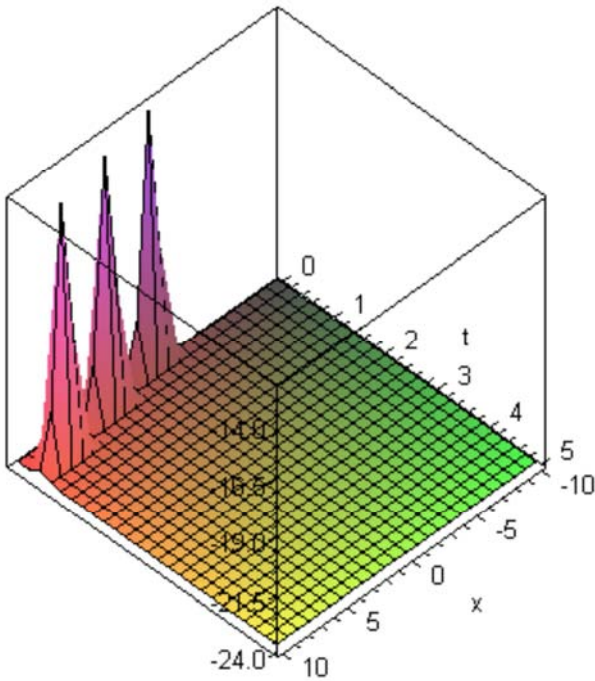


Figure 12. Soliton solution of $u_1(x,t)$ for $k_1 = 2, k_2 = 0.5, a = 2.5, b = 2.5, \lambda = 4, \mu = 2$.

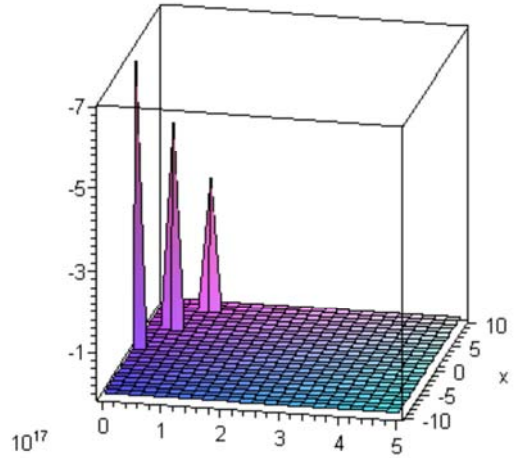


Figure 14. Soliton solution of $u_2(x,t)$ for $a = 1.5, b = 2, \lambda = 3, \mu = 1$.

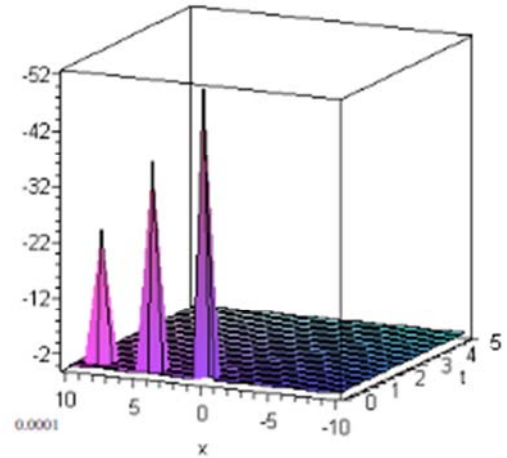


Figure 15. Soliton solution of $u_2(x,t)$ for $a = 2.5, b = 2.5, \lambda = 4, \mu = 2$.

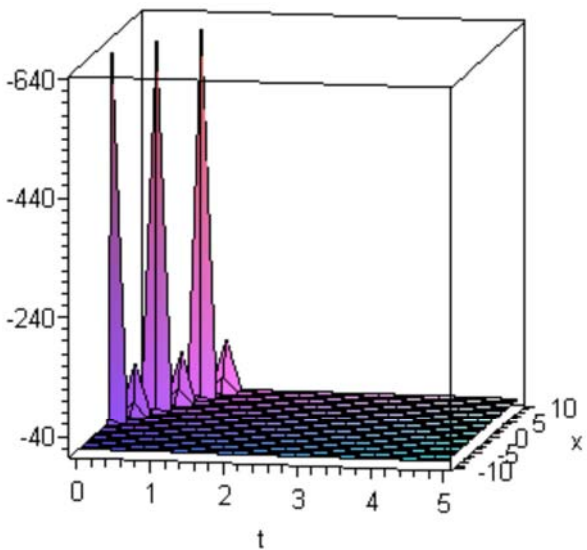


Figure 13. Soliton solution of $u_1(x,t)$ for $k_1 = 1, k_2 = 1.5, a = 1.5, b = 2, \lambda = 3, \mu = 1$.

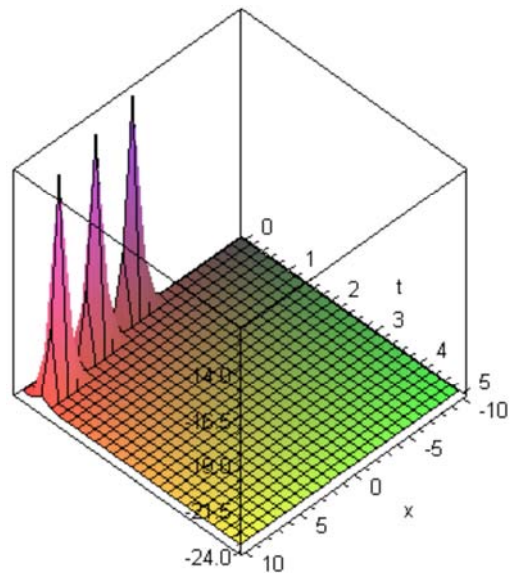


Figure 16. Soliton solution of $u_3(x,t)$ for $a = 2.5, b = 2.5, \lambda = 4, \mu = 2$.

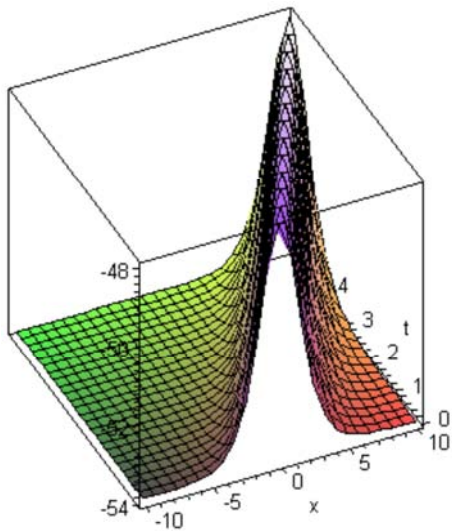


Figure 17. Soliton solution of $u_3(x,t)$ for $a = 0.5, b = 2, \lambda = 3, \mu = 2$.

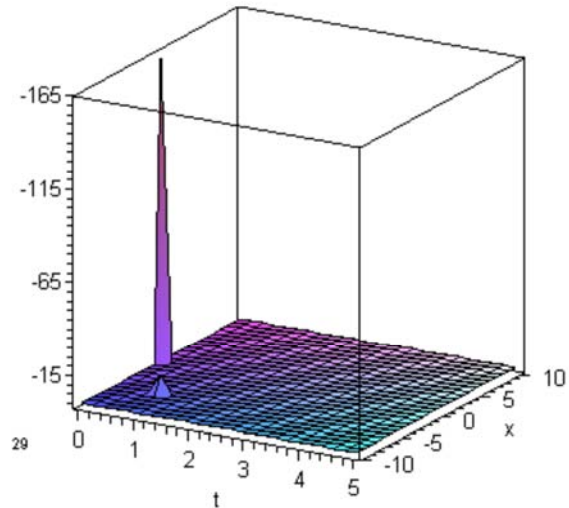


Figure 20. Soliton solution of $u_5(x,t)$ for $a = 7.5, b = 1, \lambda = 1, \mu = 2$.

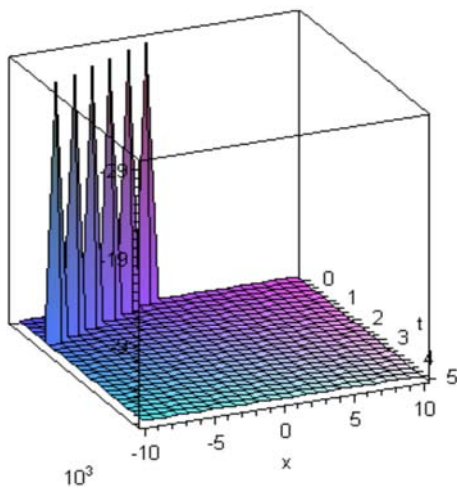


Figure 18. Soliton solution of $u_4(x,t)$ for $k_1 = 1, k_2 = 1.5, a = 0.5, b = 2, \lambda = 2, \mu = 2$.

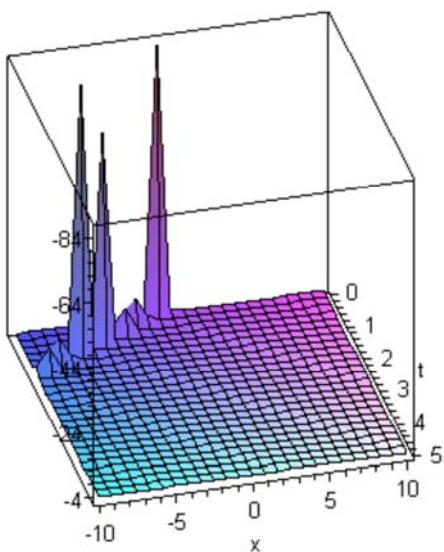


Figure 19. Soliton solution of $u_4(x,t)$ for $k_1 = 1.5, k_2 = 2, a = 7.5, b = 1, \lambda = 1, \mu = 2$.

4. Results and Discussion

From graphical representations, we note that soliton is a wave which preserves its shape after it has collided with another wave of the same kind. By solving nonlinear evolution equations included foam drainage equation and fourth order evolution equation, we attain desired solitary wave solutions for different values of random parameters. The solitary wave moves toward right if the velocity is positive. It turns in left directions if the velocity is negative. The amplitudes and velocities are controlled by various parameters. Figures signify graphical representation for different values of parameters. Figure 1 to 11 represents periodic wave solution for different values of parameters γ, λ and μ . The soliton solutions that are shown in figure 12 to 15 and figure 17 to 20 represents solitary wave solutions for different values of parameters λ and μ . Figure 16 shows peakons solution by using values of parameters as $\lambda = 4, \mu = 2$. In all cases, for various values of parameters, we attain identical solitary wave solutions which obviously show that the final solution is not effectively based upon these parameters. So, we can choose arbitrary values of such parameters as input to our simulations.

5. Conclusion

In this paper by use of (G'/G) -expansion method, the more general and new exact solutions of NLEEs has been obtained. For this, nonlinear foam drainage equation and fourth order evolution equation are considered. We attain desired soliton solutions of various types for different values of parameters. It is guaranteed the accuracy of the attain results by backward substitution into the original equation with Maple software. The scheming procedure of this method is simplest, straight and productive. It is observed that the under study technique is more reliable and have minimum computational task, so widely applicable. In precise one can say this method is quite competent and much operative for evaluating exact solution of NLEEs. The validity of given algorithm is totally hold up with the help of the

computational work, and the graphical representations. Results obtained by this method are very encouraging and reliable for solving any other type of NLEEs. The graphical representations clearly indicate the solitary solutions.

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